Ab initio Molecular Dynamics Studies of Anomalous Proton Transport Mechanisms in Various Hydrogen-bonded Media

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### Talk Outline

Proton transport in ammonium perchlorate crystal

Proton transport in water and acidic solutions

Proton transport in liquid methanol

Proton transport in methanol/water mixtures

#### SEE POSTER:

# Proton transfer in water and at the liquid-vapor interface: a combined ab initio and polarizable multi-state EVB molecular dynamics approach

G. Brancato, H. -S. Lee and MET



#### Ab initio molecular dynamics

R. Car and M. Parrinello, Phys. Rev. Lett. 55, 2471 (1985). Given a system of N nuclei with positions  $\mathbf{R}_1, ..., \mathbf{R}_N \equiv \mathbf{R}$ 

Classical evolution:

 $M_I \ddot{\mathbf{R}}_I = \mathbf{F}_I(\mathbf{R})$ 

 $= -\langle \Psi_0(\mathbf{R}) | \nabla_I H_{\text{elec}}(\mathbf{R}) | \Psi_0(\mathbf{R}) \rangle - \nabla_I U(\mathbf{R})$ 

 $\frac{\text{Kohn-Sham Density Functional Theory}}{E[\{\psi_i\}, \{\mathbf{R}\}]} = -\frac{1}{2} \sum_i \langle \psi_i | \nabla^2 | \psi_i \rangle + \frac{1}{2} \int d\mathbf{r} \, d\mathbf{r}' \, \frac{n(\mathbf{r}) n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \\ + E_{\text{xc}}[n] + \sum_{I=1}^N \int d\mathbf{r} \, n(\mathbf{r}) V_{\text{ext}}(|\mathbf{r} - \mathbf{R}_I|) + U(\mathbf{R}) \\ n(\mathbf{r}) = \sum_i |\psi_i(\mathbf{r})|^2 \qquad \langle \psi_i | \psi_j \rangle = \delta_{ij}$ 

Car-Parrinello dynamics:

$$\mu |\ddot{\psi}_i\rangle = -\frac{\partial E}{\partial \langle \psi_i |} + \sum_{i,j} \Lambda_{ij} |\psi_j\rangle$$
$$M_I \ddot{\mathbf{R}}_I = -\nabla_I E$$

## **Ammonium perchlorate crystal**

- Thermally labile materials used as an oxidizer in rocket fuels.
- Undergoes an orthorhombic -> cubic structural phase transition at *T* = 513 K.
- In the cubic phase, experiments find proton conductivity increases by factor of 5 upon doping with neutral ammonia.
- Thought that rotational dynamics of the ions plays a role in the conduction mechanism.



Pure Cubic AP Crystal

Cubic AP crystal + ammonia



L. Rosso and MET, Solid State Ionics 161, 219 (2003)

# **Simulation specifics**

L. Rosso and MET, Solid State Ionics 161, 219 (2003)

- Temperature: Orthorhombic 300 K, Cubic 530 K
- Cubic: 15.26 x 15.26 x 7.63 Å
- Ortho: 9.2 x 11.6 x 7.63 Å
- Simulation Lengths: Pure: 15 ps, Doped 30 ps
- BLYP functional
- Martins-Troullier pseudopotentials

# **Calculation of infrared spectrum**

Infrared spectrum from autocorrelation function of total dipole moment in the harmonic approximation:

$$\alpha(\omega) = \frac{4\pi\beta\omega^2}{3\hbar n(\omega)cV\varepsilon_0} \int d\omega \, e^{-i\omega t} \left\langle \mathbf{M}(0)\cdot\mathbf{M}(t)\right\rangle_{cl}$$
$$= \frac{4\pi\beta}{3\hbar n(\omega)cV\varepsilon_0} \int d\omega \, e^{-i\omega t} \left\langle \dot{\mathbf{M}}(0)\cdot\dot{\mathbf{M}}(t)\right\rangle_{cl}$$

Total electronic position operator:  $\mathbf{r} = \sum_{i} \mathbf{r}_{i}$ 

Electronic contribution: [R. Resta, Phys. Rev. Lett. 80, 1800 (1998)]

$$M_{\alpha}^{(\text{elec})} = \frac{eL}{2\pi} \operatorname{Im} \ln \det \mathbf{R}_{\alpha} \overline{\phantom{a}} e^{-2\pi i r_{\alpha}/L} \Psi_{0}(\tau)$$

$$R_{\alpha,ij} = \left\langle \psi_i \right| e^{-2\pi i r_a / L} \left| \psi_j \right\rangle$$













(Ъ)







(a)







NH<sub>4</sub>+ D (doped)/D(pure)=4.9 σ(doped)/σ(pure)=4.5 H. Wise, JPC **71**, 2843 (1967)

 $CIO_4^-$ 

### **Rotational mean-square displacements**



### **Structures of the excess proton in water**

Hydronium ion:



## **The Grotthuss mechanism in water**

MET, *et al*, JPC, **99**, 5749 (1995); JCP **103**, 150 (1995)
D. Marx, MET, M. Parrinello, *Nature* **397**, 601 (1999).
N. Agmon, *Chem. Phys. Lett.* **244**, 456 (1995)
T. J. F. Day, *et al. J. Am. Chem. Soc.* **122**, 12027 (2000)



Solvent coordinate view:

P. M. Kiefer, J. T. Hynes *J. Phys. Chem. A* **108**, 11793 (2004)

### **The Grotthuss mechanism in water**

Second solvation shell H-bond breaking followed by formation of intermediate Zundel complex:



## **The Grotthuss mechanism in water**

Transfer of proton resulting in ``diffusion'' of solvation structure:



## **Decomposition of the dipole moment**

R. Iftimie and MET, J. Chem. Phys. (in press)

Dipole moment:

$$M_{\alpha}^{(\text{elec})} = \frac{eL}{2\pi} \text{Im ln det R}_{\alpha}$$

If **R** is diagonal with eigenvalues **r**:

$$\mathbf{M} = \sum_{i} \boldsymbol{\mu}_{i} = \sum_{m} \boldsymbol{\mu}_{m}$$

where

$$\boldsymbol{\mu}_m = \sum_{j \in m} \left[ Z_j \mathbf{R}_j - \sum_{k \in j} \mathbf{r}_k \right]$$

Correlation function:

$$C(t) = \left\langle \dot{\mathbf{M}}(0) \cdot \dot{\mathbf{M}}(t) \right\rangle_{cl} = \sum_{m} \left\langle \dot{\mathbf{\mu}}_{m}(0) \cdot \dot{\mathbf{M}}(t) \right\rangle_{cl}$$

### **Unitary (Gauge) Invariance and localized orbitals**

#### Total energy invariant under static unitary transformations:

$$\tilde{\psi}_i(\mathbf{r},t) = \sum_j U_{ij} \psi_j(\mathbf{r},t) = \sum_j (e^{i\mathbf{B}})_{ij} \psi_j(\mathbf{r},t)$$



If we only care about distance from center at the potential minimum, then a point  $Z=R_{min} e^{i\theta}$  is acceptable for any choice of  $\theta$ 

### **Decomposition of the dipole moment**

Dipole moment determined by the matrix

$$R_{\alpha,ij} = \left\langle \psi_i \right| e^{2\pi i r_\alpha / L} \left| \psi_j \right\rangle$$

Different components do not commute with each other

$$\begin{bmatrix} \mathbf{R}_{\alpha}, \mathbf{R}_{\beta} \end{bmatrix} \neq 0 \qquad \text{if} \quad \alpha \neq \beta$$

Cannot simulataneously diagonalize all three components. Define

$$\Omega[\psi] = \sum_{i} \left[ \left| R_{x,ii} \right|^{2} + \left| R_{y,ii} \right|^{2} + \left| R_{y,ii} \right|^{2} \right]$$

Choose *U* such that  $\Omega[\tilde{\psi}]$  is maximal. The orbitals that result are known as Wannier functions. Because they maximize the diagonal elements of the translationally invariant position operator, they spatial spread Is minimal.

DFT = BLYP, Cutoff = 80 Ry, System = 64 molecules + 1 HCl or 1  $H_2F_2$  time = 20 ps



R. Iftimie and MET, J. Chem. Phys. (in press); *ibid*, PNAS (submitted); *ibid*, JACS (in prep)

Expt. (3.6 M HBr)
Expt. (8.8 M HBr)
Theor. (0.9 M HCl)

**–** – **–** Theor. (0.9 M  $H_2F_2$ )

Theor. (0.9 M  $H_2F_2$ ) – FHF<sup>-</sup> spectrum



#### Ab Initio Path Integrals

M. E. Tuckerman, et al, JCP 104, 5579 (1996).

#### Quantum canonical partition function

$$Q = \operatorname{Tr}\left(e^{-\beta H}\right)$$
  
=  $\oint \mathcal{D}\mathbf{R}_{1}\cdots\mathcal{D}\mathbf{R}_{N}\exp\left\{-\int_{0}^{\beta}d\tau\left[\sum_{I=1}^{N}\frac{1}{2}M_{I}\dot{\mathbf{R}}_{I}^{2}(\tau)+E_{0}(\{\mathbf{R}_{I}(\tau)\})\right]\right\}$ 

Limit of discretized integral

$$Q = \lim_{P \to \infty} \left[ \prod_{I=1}^{N} \mathcal{N} \int d\mathbf{R}_{I}^{(1)} \cdots d\mathbf{R}_{I}^{(P)} \int d\mathbf{P}_{I}^{(1)} \cdots d\mathbf{P}_{I}^{(P)} \right] \\ \times \exp \left\{ -\beta \sum_{s=1}^{P} \left[ \sum_{I=1}^{N} \left( \frac{\left(\mathbf{P}_{I}^{(s)}\right)^{2}}{2M_{I}'} + \frac{1}{2\beta^{2}} M_{I} P(\mathbf{R}_{I}^{(s)} - \mathbf{R}_{I}^{(s+1)})^{2} \right) + \frac{1}{P} E_{0}(\{\mathbf{R}_{I}\}^{(s)}) \right] \right\}$$



Exact when:

$$P \rightarrow \infty$$

# **Simulation specifics**

Marx, MET, Hutter, Parrinello, Nature 397, 601 (1999).

- **Temperature:** T = 300 K
- **Box:** L = 9.87 Å
- System: 31  $H_2O + 1 H_3O^+$
- **Path integral discretization:** P = 8, Classical: P=1
- Simulation Lengths:

Quantum:322,000 configurations (40 ps)Classical (P=1):345,000 configurations (60 ps)

## **Coordinates that follow the proton**



#### **Project out the Defect Site**



## **Probability distribution functions**



### **Proton Rattling: Free Energy Profile**

For analysis: 1D proton transfer coordinate <sup>1</sup>



Proton rattling in  $H_5O_2^+$  complex:  $H_2O \stackrel{\leftrightarrow}{\cdots} H^* \stackrel{\leftrightarrow}{\cdots} OH_2$ 

### **Quantum delocalization of structural defect**

D. Marx, MET, J. Hutter and M. Parrinello Nature 397, 601 (1999)



## **Proposed Structural Diffusion Mechanism in MeOH**



H. -C. Chang, et al, J. Phys. Chem. A 103, 2941 (1999).

## **Bond Cutting**



Use of monovalent Carbon pseudopotential

U. Röthlisberger, private comm.

Y. Zhang, T. Lee, W. Yang JCP 110, 46 (1999).



## **Results on structure of liquid methanol**

DFT Level: **BLYP** Force Field: AMBER95 CO bond cutting: Pseudobond System size: 32 molecules Sim. Length: 10 ps

Electron-methyl H pseudopotential:

$$V_{e-M}(r) = Ae^{-\gamma r}$$

A=18.5763 kcal/mol

γ = 0.5292 Å



J. A. Morrone and MET, Chem. Phys. Lett. 370, 406 (2003)



A. Adya, et al. J. Chem. Phys. 112, 4231 (2000)

### **H-bond chain length statistics**





# **Simulation specifics**

J. A. Morrone and MET, *J. Chem. Phys.* **117**, 4403 (2002). J. A. Morrone and MET, *Chem. Phys. Lett.* **370**, 406 (2003).

- **Temperature:** T = 300 K
- **Box:** L = 12.93 Å
- System: 32 MeOH and 32 MeOH + 1 H<sup>+</sup>
- All H are D
- Simulation Lengths: 60 ps
- BLYP treatment of OH group
- AMBER treatment of Me
- Monovalent pseudopotential in "cut" bond
- e-M pseudopotential  $V(r) = Ae^{-\gamma r}$

## **Geometry of the protonated methanol dimer**





# **The Hazard Plot**

Suppose a system which can undergo a particular process has survived up to time t without undergoing the process. Let h(t)dt be the probability that the process will occur between t and t + dt (called the hazard). Define the cumulative hazard:

$$H(t) = \int_0^t dt' h(t')$$

Plot H(t) vs. transition time.

If process is random, then at long time,  $h(t) = \lambda$ ,  $H(t) = \lambda t$ .

E. Helfand, J. Chem. Phys. 69, 1010 (1978).

# **Proton transfer event statistics**



# **Defect H-bond length distributions**





## **Methanol-water mixtures**

Competition between chain formation of methanol and tetrahedral network of water

Recent neutron diffraction experiments suggest that between X=0.27 and x=0.54, molecular segregation occurs, leading to separate "percolating networks".

S. Dixit, et al. Nature 416, 829 (2002); L. Dougan, et al. J. Chem. Phys. 121, 6456 (2004)

Water cluster sizes mostly less than 10 at x=0.27 and less than 20 for x=0.54.

Suppression of methanol tendency to form hydrogen-bonded chains as *x* increases.

### From a classical MD run using 1024 MeOH + 1024 $H_2O$ (x=0.5)



# **Simulation specifics**

K. E. Haslinger, J. A. Morrone and MET (in preparation)

- **Temperature:** T = 300 K
- **Box:** L = 11.46 Å (x=0.5), 10.69 Å (x=0.25)
- System: 16 MeOH and 16 H<sub>2</sub>O (+ 1 H<sup>+</sup>): x=0.5
- System: 8 MeOH and 24  $H_2O$ : x=0.25
- All H are D
- Simulation Lengths: 100 ps
- Vanderbilt Ultrasoft Pseudopotentials

### **Water-water radial distribution functions**



### **Methanol-methanol radial distribution function**



#### From ab initio MD run (x=0.5)



### A simple kinetic model

Let A be the most active hydrogen bond. Let B be an adjacent hydrogen bond. Let C be a more "distant" situation.

$$A \underset{k_{-1}}{\overset{k_{1}}{\underset{k_{-1}}{\longrightarrow}}} B$$
$$B \underset{k_{2}}{\longrightarrow} C$$

$$\frac{d[A]}{dt} = -k_1[A] + k_{-1}[B]$$
$$\frac{d[B]}{dt} = -(k_2 + k_{-1})[B] + k_1[A]$$

### Solution of the kinetic model

$$[A](t) = \frac{1}{2\lambda} \Big[ (\lambda - k_1 + K) e^{-(k_1 + K - \lambda)t/2} + (\lambda + k_1 - K) e^{-(k_1 + K + \lambda)t/2} \Big]$$
  

$$= [A]_{slow}(t) + [A]_{fast}(t)$$
  

$$[B](t) = \frac{k_1}{\lambda} \Big[ e^{-(k_1 + K - \lambda)t/2} - e^{-(k_1 + K + \lambda)t/2} \Big]$$
  

$$= [B]_{slow}(t) + [B]_{fast}(t)$$
  

$$\lambda = \sqrt{(k_1 - K)^2 + 4k_1k_{-1}}$$
  

$$K = k_2 + k_{-1}$$

### Using the model with the MD data





Correlation functions for a proton in a 50:50 methanol/water mixture

From a LLSQ fit:

 $k_1 = 20.9 \text{ ps}^{-1}$   $k_{-1} = 22.6 \text{ ps}^{-1}$   $k_2 = 0.8 \text{ ps}^{-1}$ 

#### Correlation functions for hydroxide mobility in water



### **Calculated approximate branching ratios**

Reaction	<b>Fraction</b>
$H_3O^+ + H_2O \rightarrow H_2O + H_3O^+$	47.0%
$H_3O^+ + MeOH \rightarrow H_2O + MeOH_2^+$	25.0%
$MeOH_2^+ + H_2O \rightarrow MeOH + H_3O^+$	25.0%
$MeOH_2^+ + MeOH \rightarrow MeOH + MeOH_2^+$	3.0%

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